

MAGNETIC DEFLAGRATION

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Magnetic clusters:

**Mn, Mn₂, Mn₃, Mn₄, [Mn₄]₂, Mn₅, Mn₆, Mn₇, Mn₈, Mn₉, Mn₁₀,
Mn₁₁, Mn₁₂, Mn₁₃, Mn₁₆, Mn₁₈, Mn₂₁, Mn₂₄, Mn₂₆, Mn₃₀, Mn₈₄**

Fe₂, Fe₃, Fe₄, Fe₅, Fe₆, Fe₇, Fe₈, Fe₁₀, Fe₁₁, Fe₁₃, Fe_{17/19}, Fe₁₉

Ni₄, Ni₅, Ni₆, Ni₈, Ni₁₂, Ni₂₁, Ni₂₄

Co₄, Co₆, Co₁₀

**Co₂Gd₂, Co₂Dy₂, Cr₁₂, CrNi₆, CrNi₂, CrCo₃, Fe₁₀Na₂, Fe₂Ni₃,
Mn₂Dy₂, Mn₂Nd₂, V₁₅, Ho, ...**

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S = 0, 1/2, 1, 3/2, 2, 5/2, 4, 9/2, 5, 33/2

Chemistry: T. Liz, 1980; Magnetism: R. Sessoli et al., 1990s

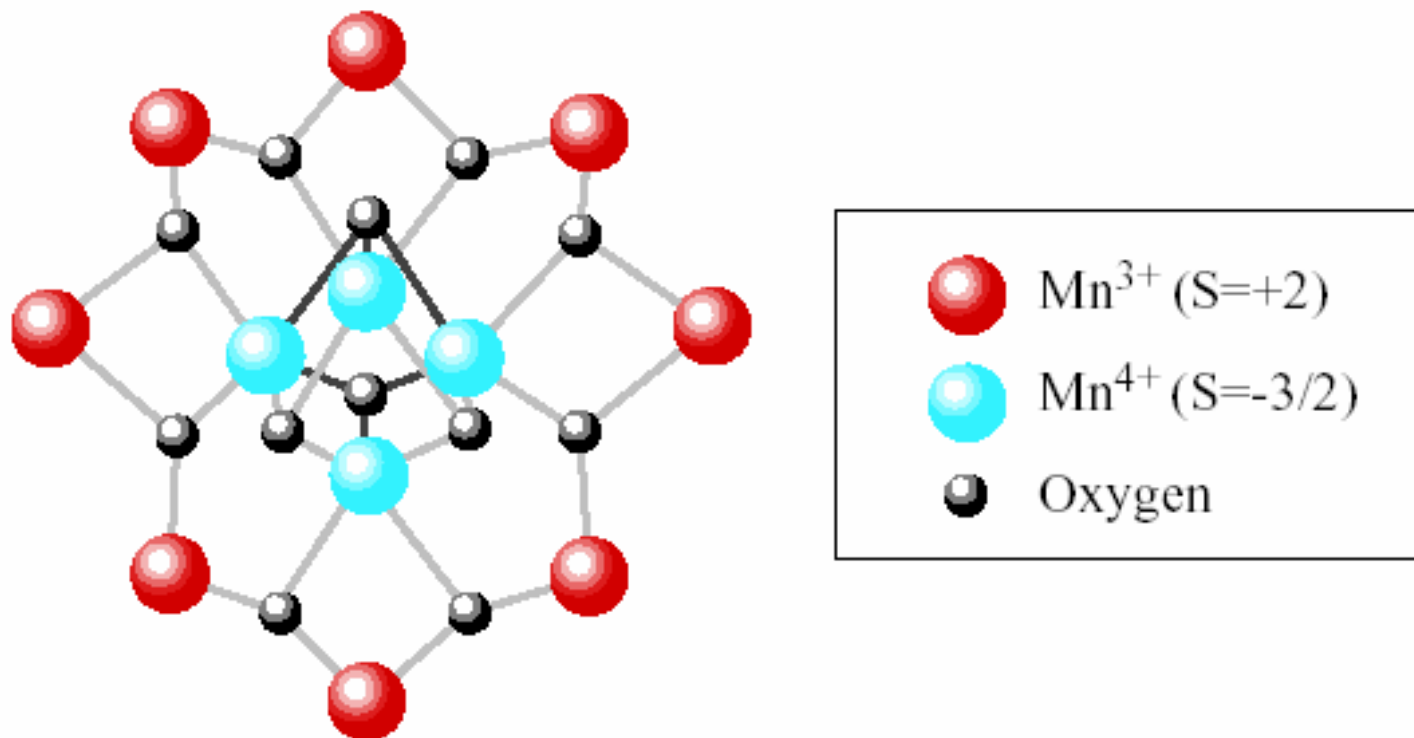
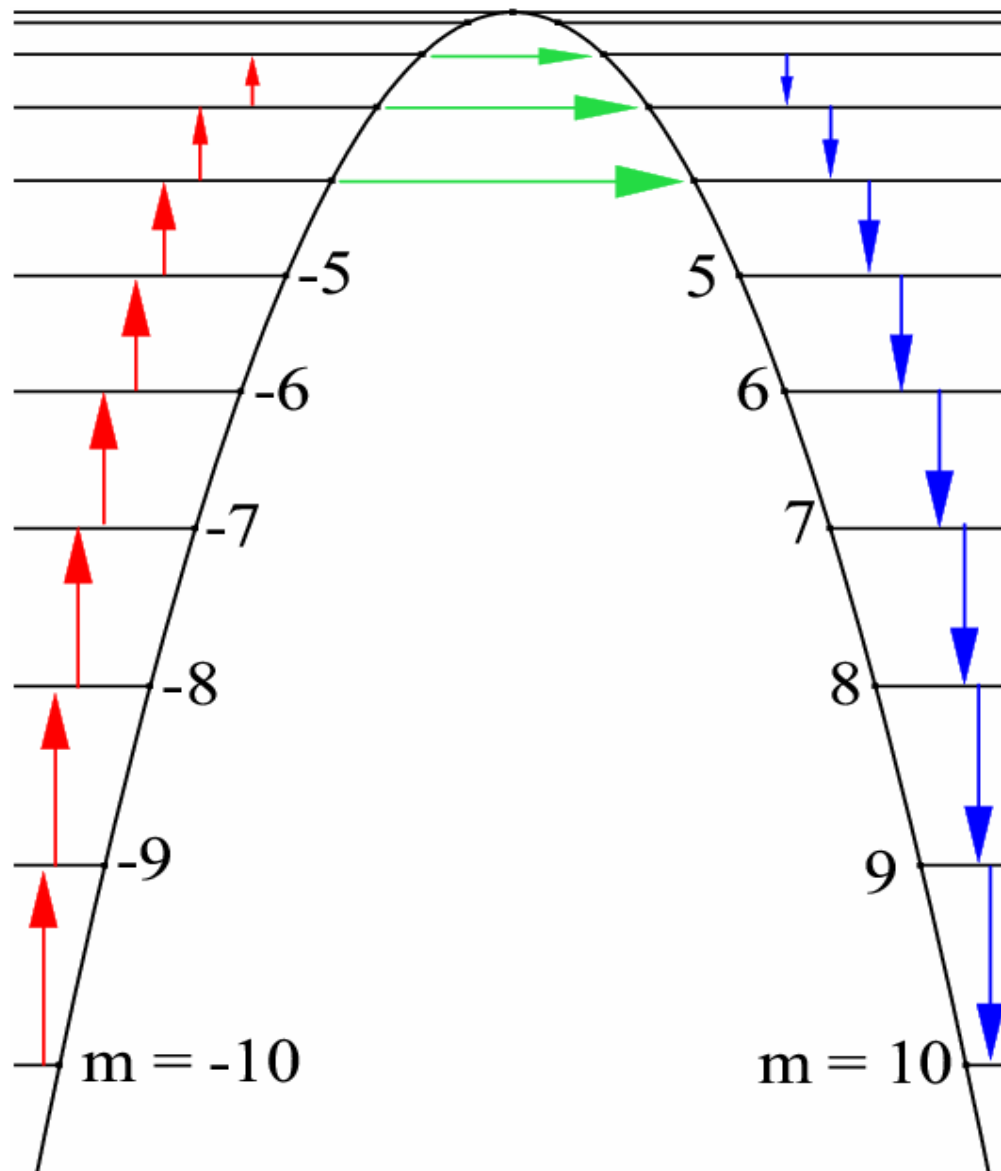


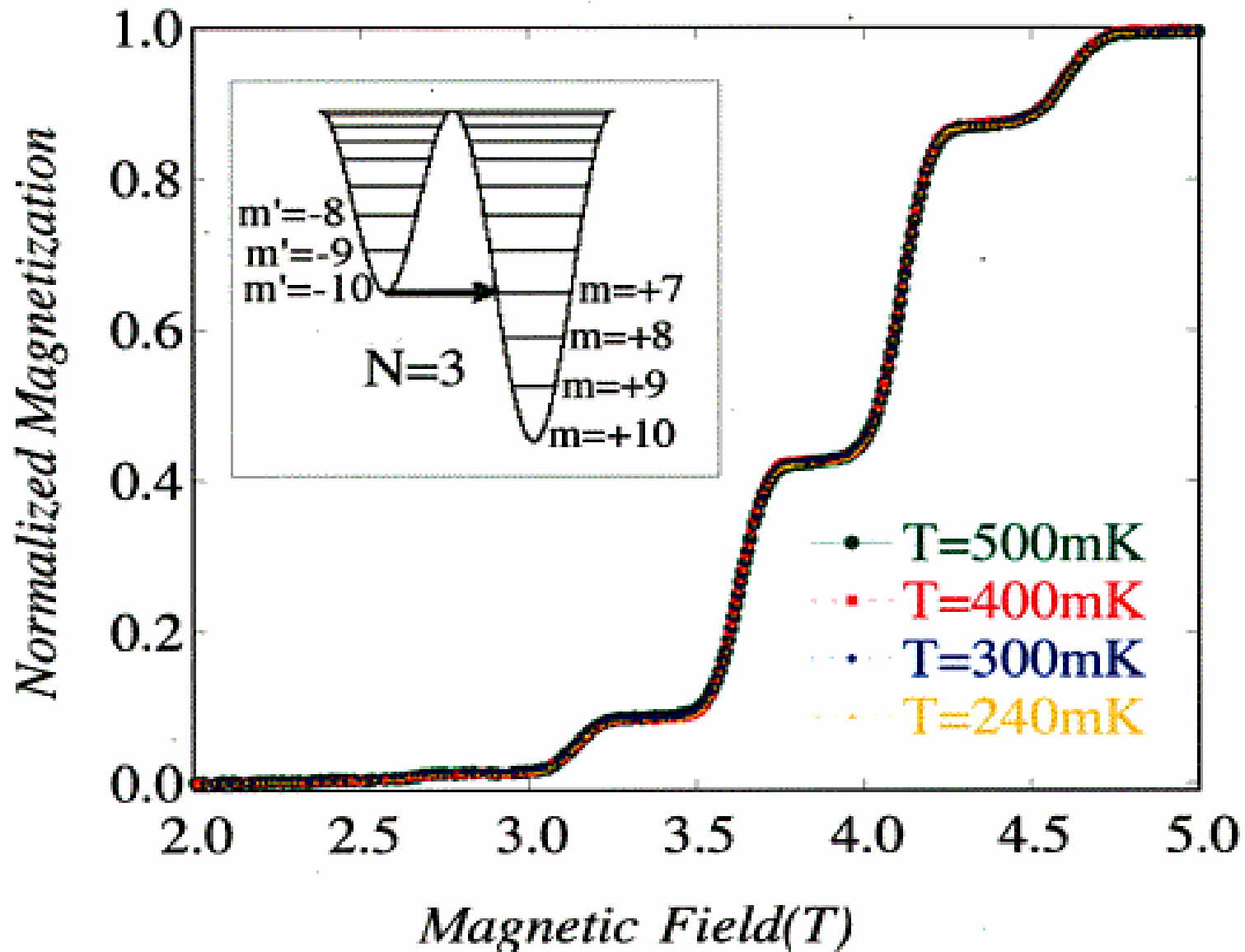
Figure 3: The magnetic core of the $S = 10$ single molecule magnet Mn_{12} -acetate. The four inner Mn ions couple anti-ferromagnetically via superexchange through oxygen bridges to the eight outer Mn ions.



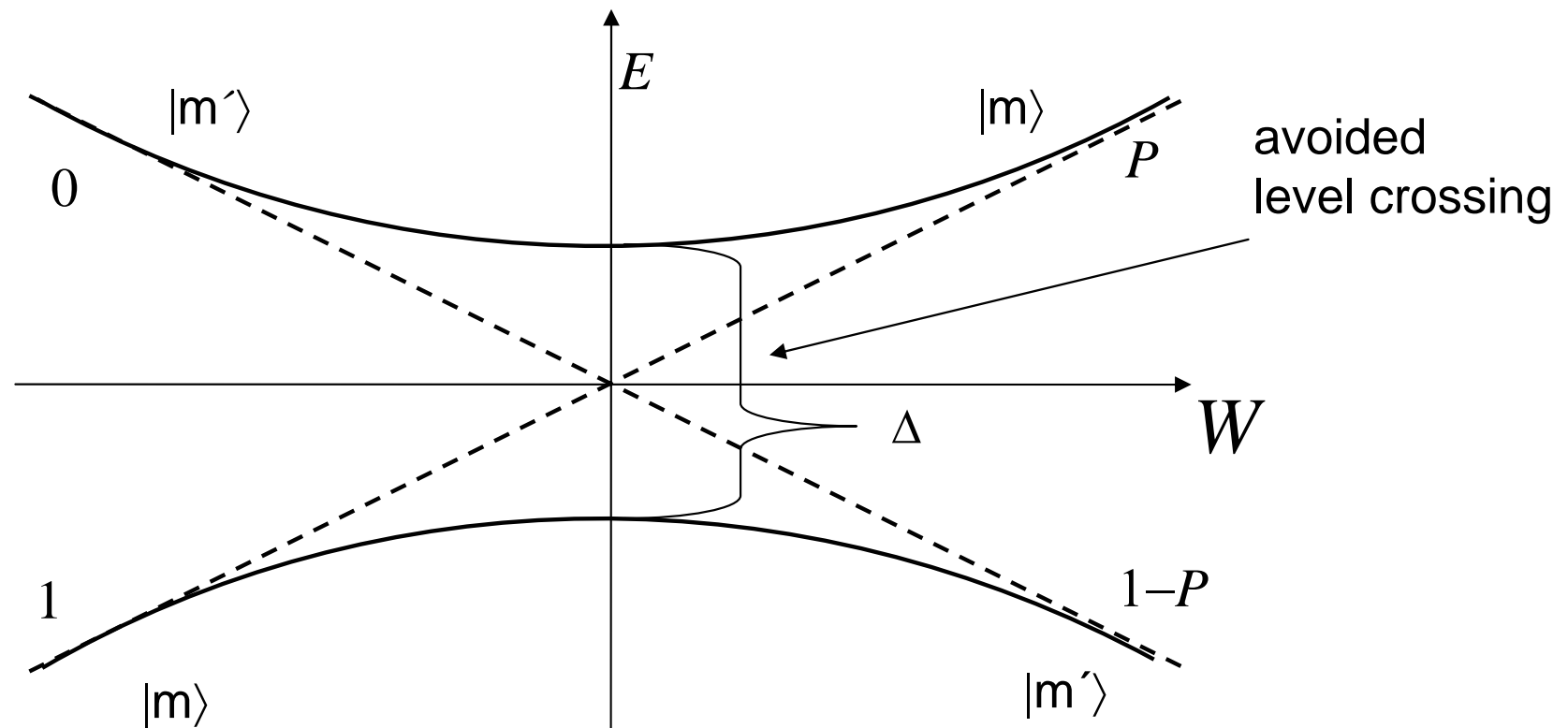
$$H_z = 0$$

Quantum Magnetization Curve

J. Friedman, M. Sarachik, J. Tejada, R. Ziolo (PRL - 1996)



Landau-Zener transitions in molecular magnets



The size of the magnetization step at each crossing of spin levels is determined by the Landau - Zener probability :

$$P = \exp\left[-\frac{\pi\Delta^2}{2\eta(dE/dt)}\right]$$

AFM dimer $[\text{Mn4}]_2$, $S=\pm 9/2$, Wernsdorfer et al., Nature-2002; Hill et al., Science-2003

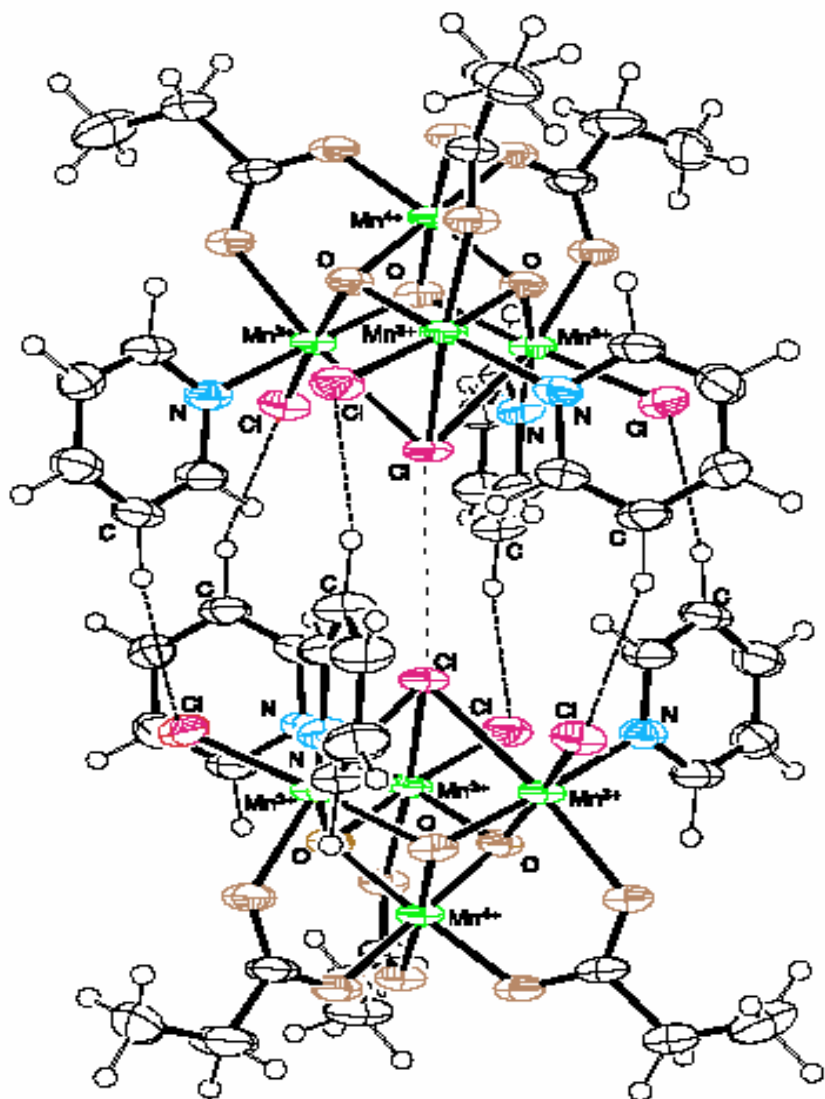


Figure 1 The structure of the $[\text{Mn}_4\text{O}_3\text{Cl}_4(\text{O}_2\text{CET})_3(\text{py})_3]_2$ dimer, denoted $[\text{Mn4}]_2$. The small circles are hydrogen atoms. The dashed lines are $\text{C}-\text{H}\cdots\text{Cl}$ hydrogen bonds and the dotted line is the close $\text{Cl}\cdots\text{Cl}$ approach. Brown, oxygen; green, manganese; red, chlorine; blue, nitrogen; py, pyridine.

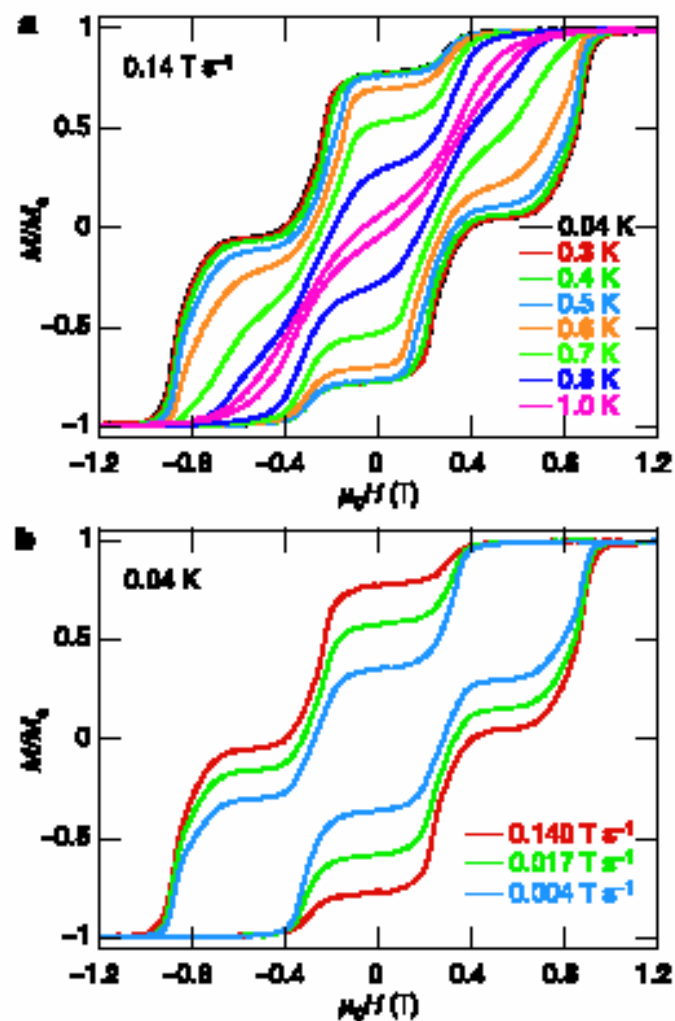
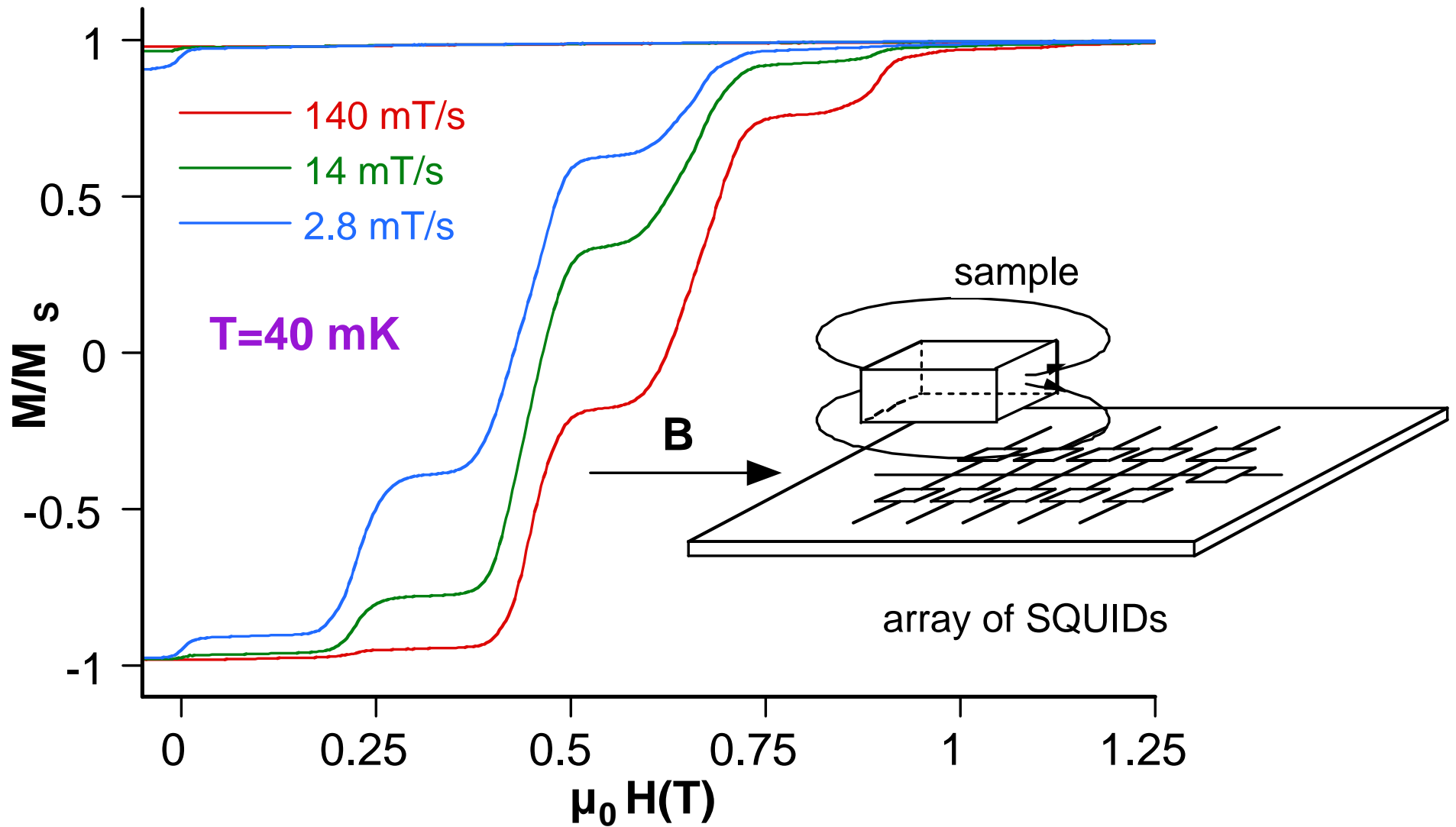


Figure 2 Magnetization (M) of $[\text{Mn4}]_2$ (plotted as fraction of maximum value M_s) versus applied magnetic field ($\mu_0 H$). **a, b**, The resulting hysteresis loops are shown at different temperatures (**a**) and different field sweep rates (**b**). We note that the loops become temperature-independent below about 0.3 K but are still sweep-rate-dependent owing to resonant quantum tunnelling between discrete energy states of the $[\text{Mn4}]_2$ dimer.

Magnetization curve of Fe-8 (W. Wernsdorfer)



Berry Phase Effects: Interference of Tunneling Paths

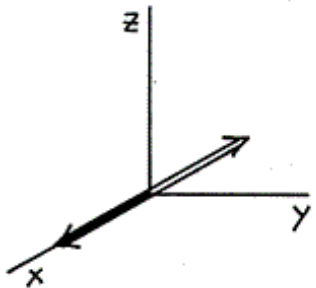
Berry phase: $A \propto e^{-iS \int d\phi (1 - \cos\theta)}$

1) $H = 0$

$$\mathcal{H} = a S_z^2 - b S_x^2$$

$$\frac{1}{2} e^{-iS\pi} + \frac{1}{2} e^{iS\pi} = \cos(S\pi)$$

[Loss, DiVincenzo, Grinstein - 1992]



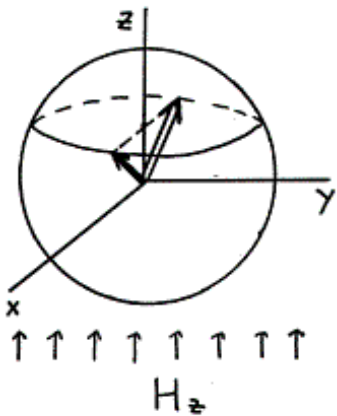
2) $H_z = 0$

$$\mathcal{H} = a S_z^2 - b S_x^2 - H_z S_z$$

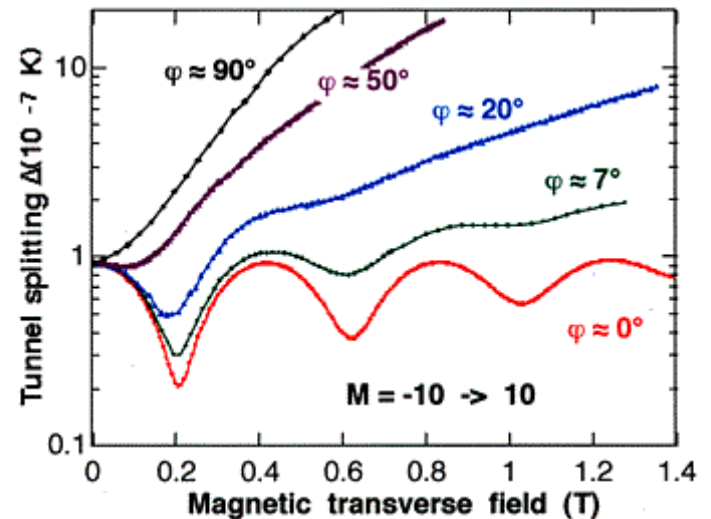
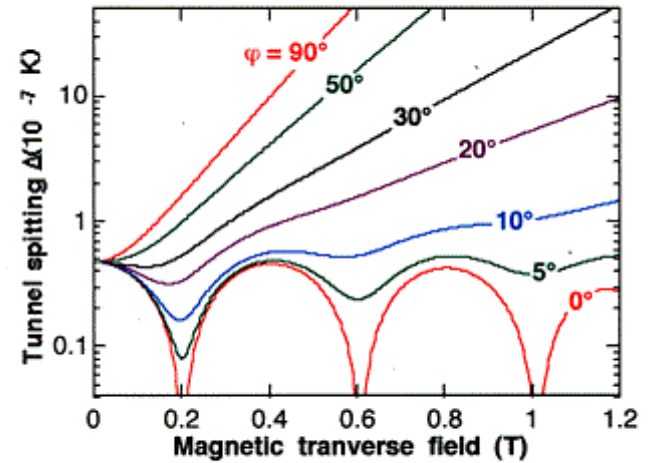
$$\frac{1}{2} e^{-iS\pi(1 - \frac{H_z}{H_c})} + \frac{1}{2} e^{iS\pi(1 - \frac{H_z}{H_c})}$$

$$= \cos[S\pi(1 - \frac{H_z}{H_c})]$$

[Garg - 1993]



Transverse field dependence of tunnel splitting

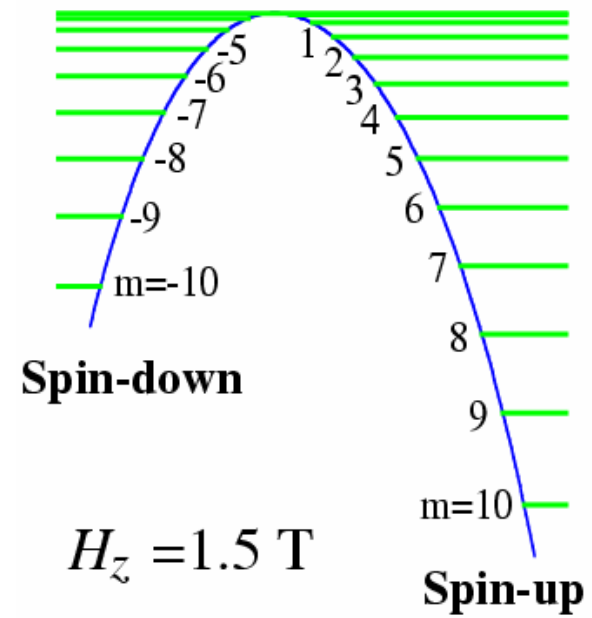
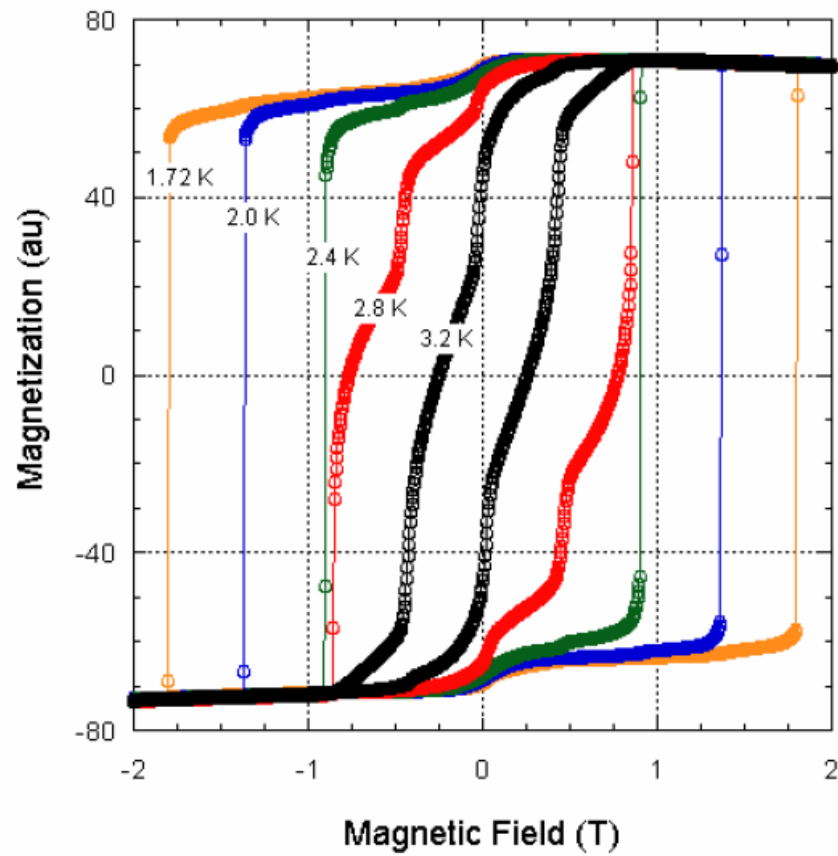


W. Wernsdorfer and R. Sessoli, *Science* 284, 133 (1999)

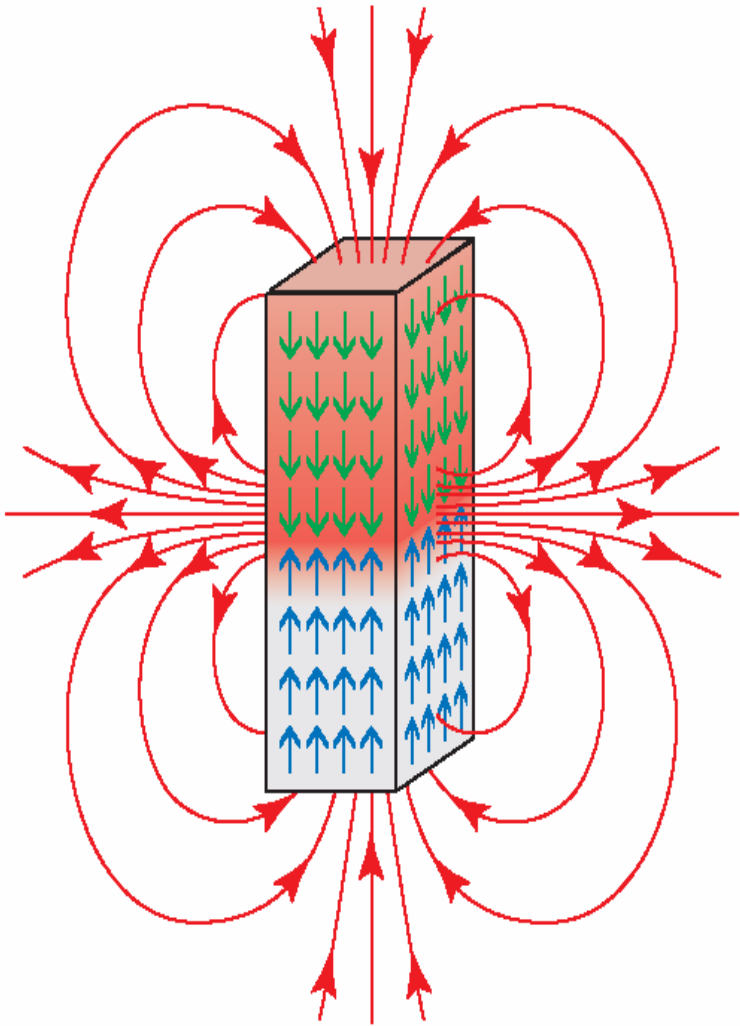
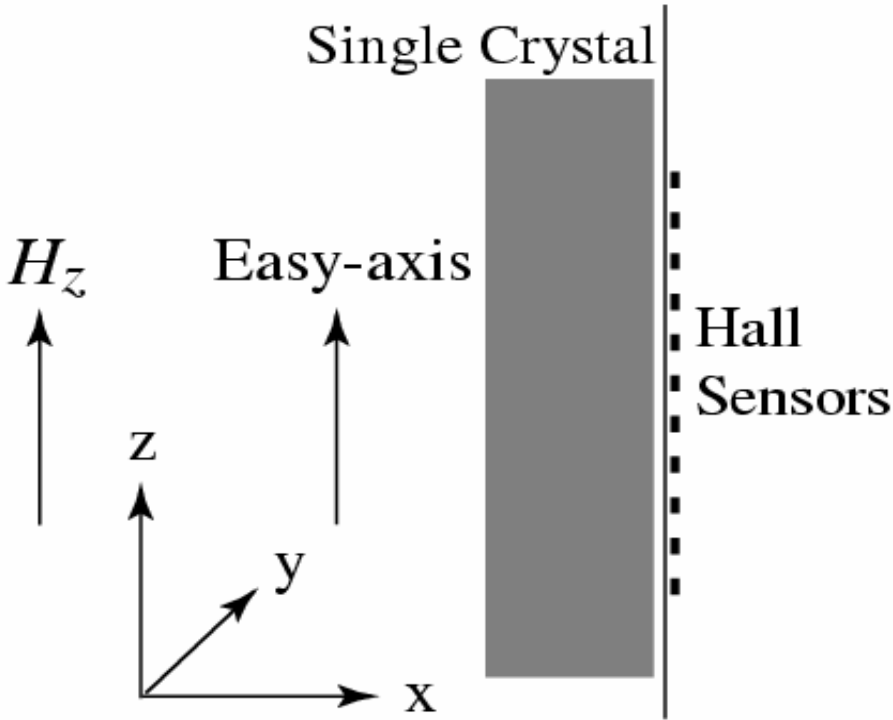
$$H = -DS_z^2 + E(S_+^2 + S_-^2) + C(S_+^4 + S_-^4) + g\mu_B \vec{S} \vec{H}$$

$$D = 0.292 \text{ K}, E = 0.046 \text{ K}, C = -2.9 \times 10^{-5} \text{ K}$$

Avalanches in Mn-12 Acetate

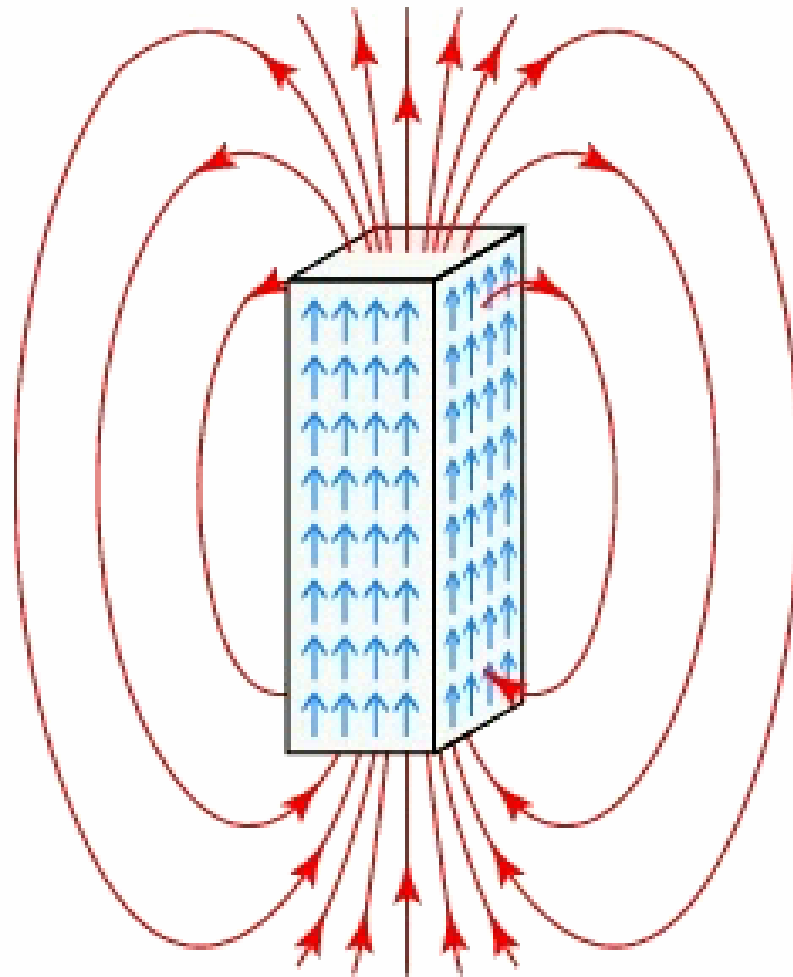


CCNY Experiment – Spring 2005, Y. Suzuki, M. Sarachik, et al.



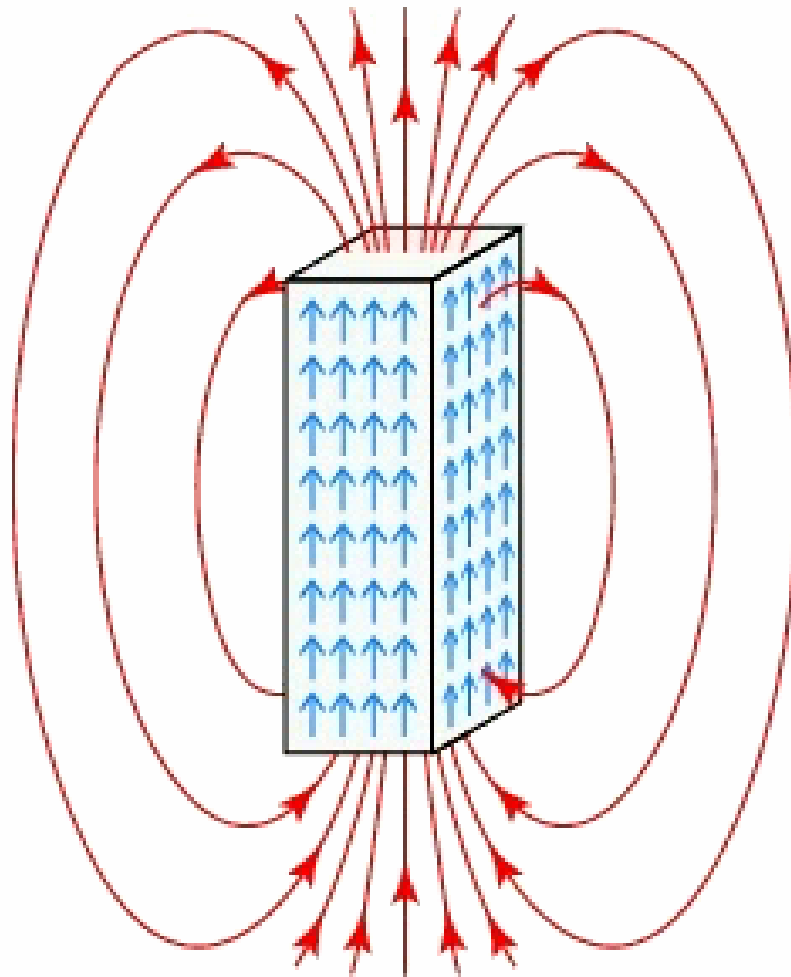
Magnetic Deflagration

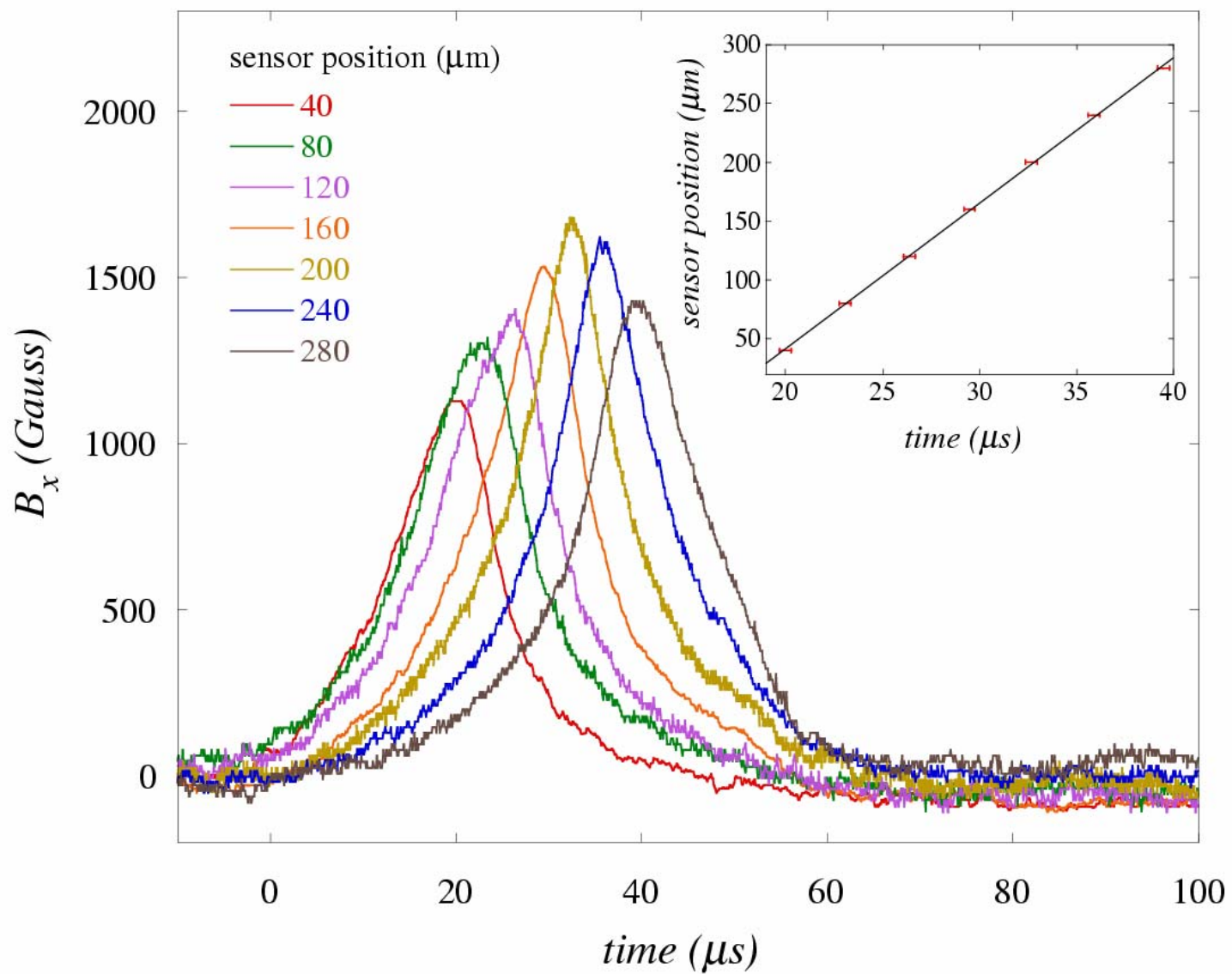
[Y. Suzuki, M. Sarachik, E. Chudnovsky, PRL – 30 Sept 2005]



Magnetic Deflagration

[Y. Suzuki, M. Sarachik, E. Chudnovsky, PRL – 30 Sept 2005]





Theory of Magnetic Deflagration

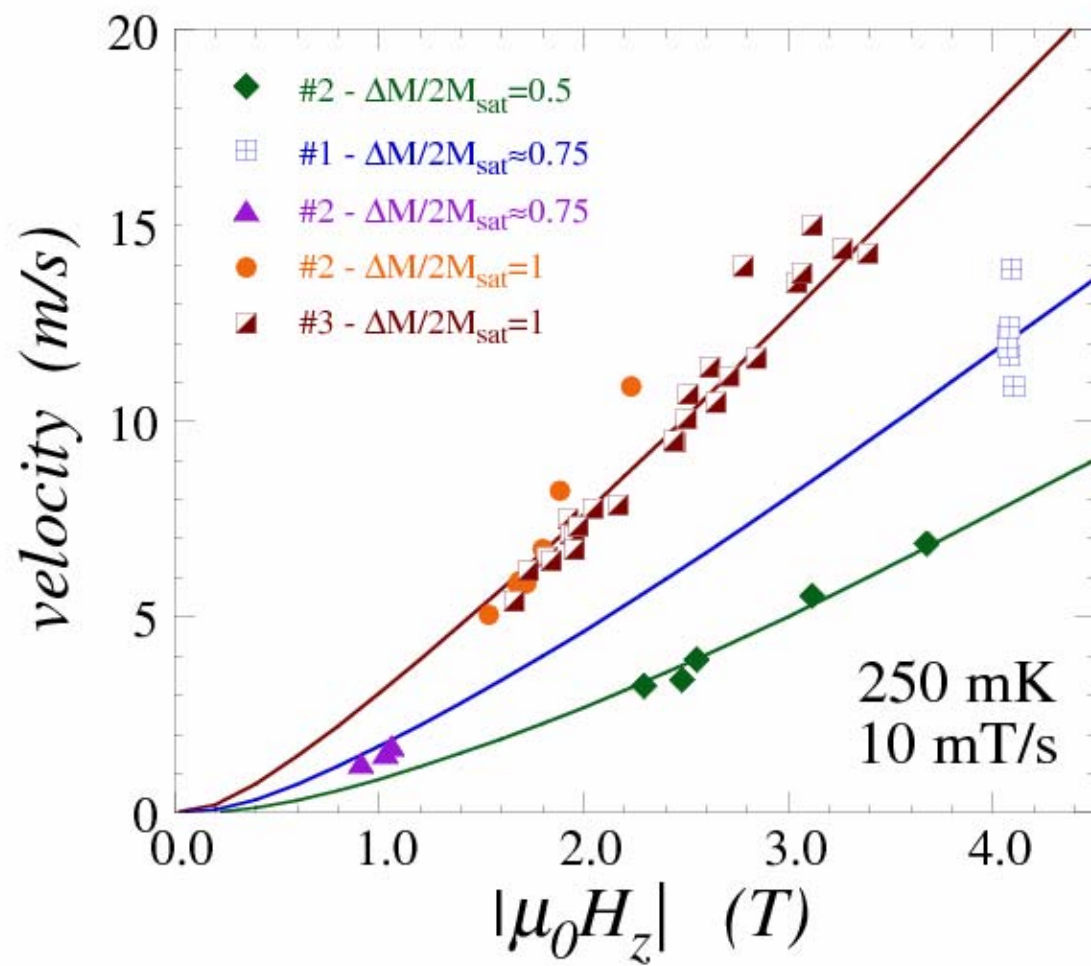
$$\frac{\partial T}{\partial t} = \kappa \nabla^2 T, \quad T = T(x - vt)$$

$$T = T_f \exp\left[-\frac{x - vt}{\delta}\right] \quad \text{at} \quad x > vt$$

$$\delta = \frac{\kappa}{v} \approx (\kappa \tau)^{1/2}$$

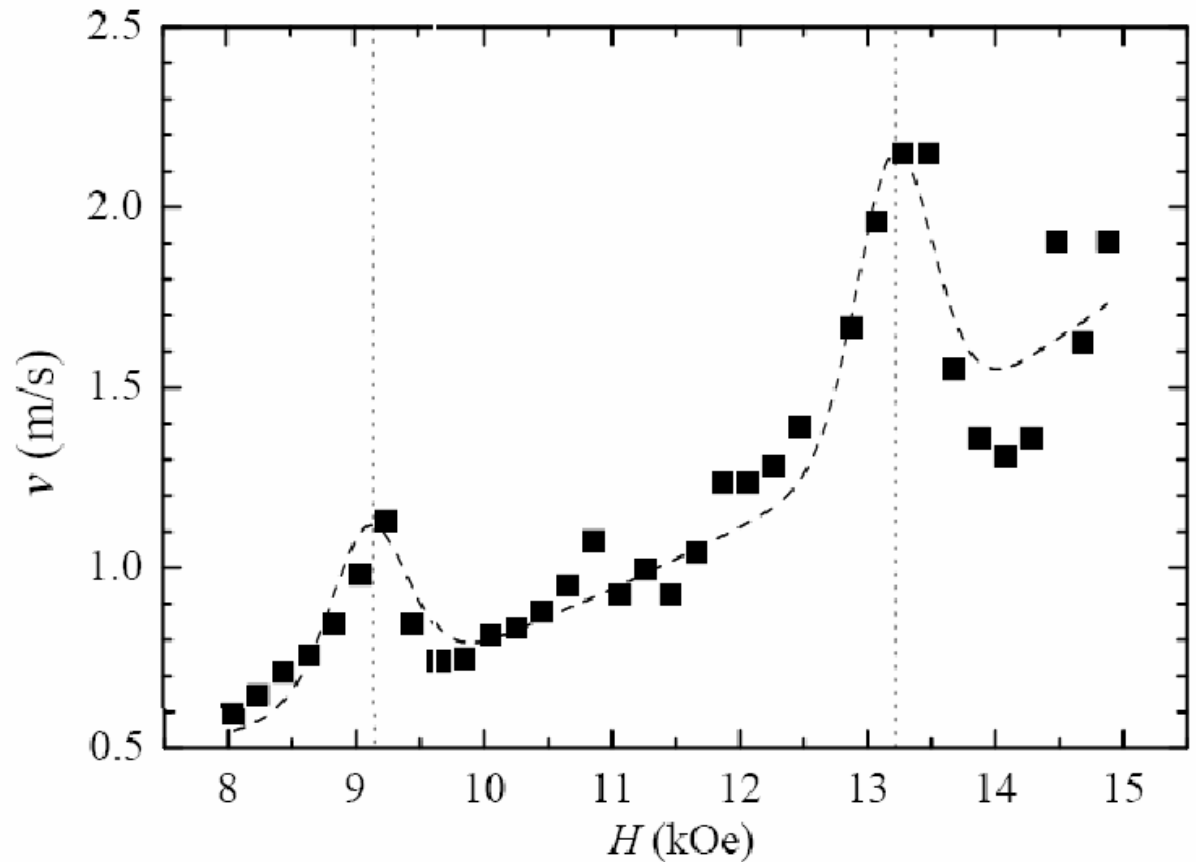
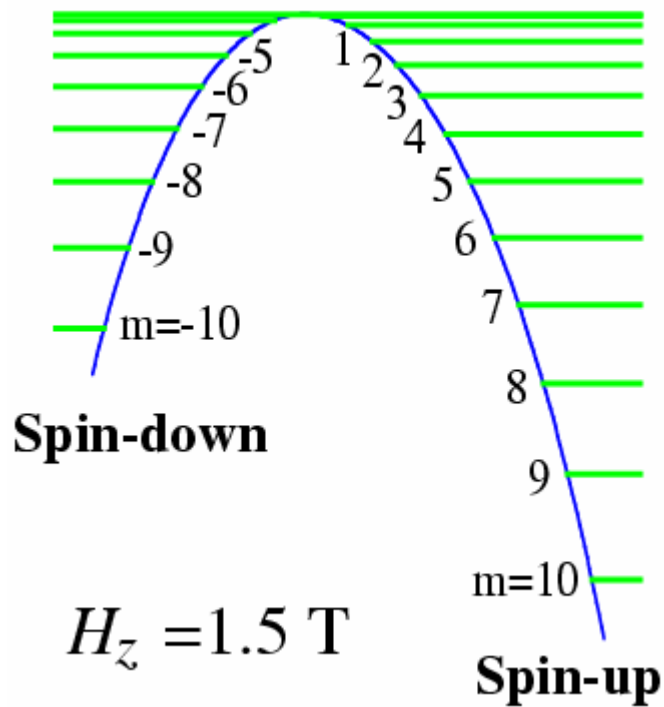
$$\tau(H) = \tau_0 \exp\left[\frac{U(H)}{k_B T}\right]$$

$$v \approx \left(\frac{\kappa}{\tau}\right)^{1/2} = \left(\frac{\kappa}{\tau_0}\right)^{1/2} \exp\left[-\frac{U(H)}{2k_B T}\right]$$



Quantum Magnetic Deflagration

[A. Hernandez-Minguez et al, PRL – Nov 2005, to appear]



CONCLUSIONS

Crystals of molecular nanomagnets exhibit two modes of magnetic relaxation:

- The slow relaxation mode, which is responsible for the staircase magnetization curve, is due to thermally assisted quantum tunneling (Landau-Zener transitions) between resonant spin levels.
- The fast magnetization reversal is equivalent to the deflagration in a flammable chemical substance. This observation opens ways of controllable non-destructive studies of deflagration and, possibly, detonation.